

# ON MULTIDIMENSIONAL ANALOGUE OF MARCHAUD FORMULA FOR FRACTIONAL RIESZ-TYPE DERIVATIVES IN DOMAINS IN $\mathbb{R}^n$

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Dedicated to Professor H.M. Srivastava, on the occasion of his 65th Birth Anniversary

### Abstract

There is given a generalization of the Marchaud formula for one-dimensional fractional derivatives on an interval (a, b),  $-\infty < a < b \leq \infty$ , to the multidimensional case of functions defined on a region in  $\mathbb{R}^n$ :

$$\mathbb{D}_{\Omega}^{\alpha}f(x) = c(\alpha) \left[ a_{\Omega}(x)f(x) + \int_{\Omega} \frac{f(x) - f(y)}{|x - y|^{n + \alpha}} dy \right], \qquad x \in \Omega, \qquad 0 < \alpha < 1,$$

which is the Riesz fractional derivative of the zero continuation of f(x) from  $\Omega$  to the whole space  $\mathbb{R}^n$ ,  $c(\alpha)$  being a certain constant. A special attention is paid to the role of the coefficient  $a_{\Omega}(x)$ , which in the multidimensional case is estimated in terms of the power of the distance of the point x to the boundary  $\partial\Omega$ . In the case when  $\Omega$  is a ball, this function is calculated explicitly in terms of the Gauss hypergeometric function.

It is also shown that the operator  $\mathbb{D}_{\Omega}^{\alpha}$  acts boundedly from the range of the Riesz potential operator  $I_{\Omega}^{\alpha}(L_p(\Omega))$  to  $L_p(\Omega), 1 .$ 

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#### 1. Introduction

The Marchaud formula

$$\mathbb{D}^{\alpha}f(x) = \frac{1}{\Gamma(1-\alpha)}\frac{f(x)}{(x-a)^{\alpha}} + \frac{\alpha}{\Gamma(1-\alpha)}\int_{0}^{x-a}\frac{f(x) - f(x-t)}{t^{1+\alpha}}\,dt, \qquad x > a,$$
(1.1)

for fractional derivatives of order  $0 < \alpha < 1$ , is well known, see [5], Subsection 13.1, which is a "difference form" of the Riemann-Liouville fractional derivative.

We introduce a multidimensional analogue of this formula in domains in  $\mathbb{R}^n$  adjusted for the Riesz fractional derivatives, see about Riesz fractional derivatives  $\equiv$  hypersingular integrals in [5], Section 26.

This generalized Marchaud formula for a domain  $\Omega$  has the form

$$\mathbb{D}_{\Omega}^{\alpha}f(x) = c(\alpha) \left[ a_{\Omega}(x)f(x) + \int_{\Omega} \frac{f(x) - f(y)}{|x - y|^{n + \alpha}} dy \right], \qquad x \in \Omega, \qquad (1.2)$$

where  $0 < \alpha < 1$  and

$$a_{\Omega}(x) = \int_{R^n \setminus \Omega} \frac{dy}{|x - y|^{n + \alpha}} \quad \text{and} \quad c(\alpha) = \frac{2^{\alpha} \Gamma\left(1 + \frac{\alpha}{2}\right) \Gamma\left(\frac{n + \alpha}{2}\right) \sin \frac{\alpha \pi}{2}}{\pi^{1 + \frac{n}{2}}}.$$

We prove a property of the function  $a_{\Omega}(x)$  important for further applications, namely, that the function  $a_{\Omega}(x)$  behaves, generally speaking as  $[\delta(x)]^{-\alpha}$ , as x approaches the boundary  $\partial\Omega$ , where  $\delta(x) = dist(x, \partial\Omega)$  is the distance of a point  $x \in \Omega$  to the boundary. We also show that in the case of the ball  $\Omega = \{x \in \mathbb{R}^n : |x| < 1\}$  the function  $a_{\Omega}(x)$  may be explicitly calculated. Finally, we show that the operator  $\mathbb{D}^{\alpha}_{\Omega}$  has some features of the operator inverse to the Riesz potential operator over  $\Omega$ .

#### 2. Definition

As is known([5], Section 26), for functions defined on the whole Euclidean space, the Riesz derivative  $\mathbb{D}^{\alpha}f = F^{-1}|\xi|^{\alpha}Ff$ , where F stands for the Fourier transforms, in the case  $0 < \alpha < 1$  has the form

$$\mathbb{D}^{\alpha} f = c(\alpha) \int_{\mathbb{R}^n} \frac{f(x) - f(x - y)}{|y|^{n + \alpha}} \, dy.$$

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Let now  $\Omega$  be a domain in  $\mathbb{R}^n$  and f(x) a function defined in  $\Omega$ . We introduce the fractional Riesz-type derivative  $\mathbb{D}_{\Omega}^{\alpha}f(x)$  of f as a restriction onto  $\Omega$  of the Riesz derivative of the zero extension of f to the whole space  $\mathbb{R}^n$ . Namely, let

$$\mathcal{E}_{\Omega}f(x) = \begin{cases} f(x), & x \in \Omega\\ 0, & x \in \mathbb{R}^n \setminus \Omega \end{cases} = : \tilde{f}(x).$$

Then, by definition

$$\mathbb{D}_{\Omega}^{\alpha}f(x) := r_{\Omega}\mathbb{D}^{\alpha}\mathcal{E}_{\Omega}f(x) = c(\alpha)\int_{\mathbb{R}^{n}} \frac{f(x) - \tilde{f}(x-y)}{|y|^{n+\alpha}} \, dy, \qquad x \in \Omega, \quad (2.1)$$

where  $r_{\Omega}$  stands for the operator of restriction onto  $\Omega$ . By splitting the integration in (2.1) to  $\int_{\Omega} + \int_{\mathbb{R}^n \setminus \Omega}$ , we easily arrive at (1.2).

In what follows, the convergence of the integral in (1.2) is interpreted as

$$\int_{\Omega} \frac{f(x) - f(y)}{|x - y|^{n + \alpha}} \, dy = \lim_{\varepsilon \to 0} \int_{\substack{y \in \Omega \\ |y - x| > \varepsilon}} \frac{f(x) - f(y)}{|x - y|^{n + \alpha}} \, dy, \qquad x \in \Omega.$$
(2.2)

Obviously, this integral absolutely converges on functions f satisfying the Hölder condition of order  $\lambda > \alpha$  in  $\overline{\Omega}$ . The integral, defining  $a_{\Omega}(x)$  is always convergent, see estimation in (3.3).

# 3. Boundary behavior

We recall that a domain  $\Omega$  is said to have the cone property, if for every  $x \in \overline{\Omega}$  there exists a finite cone  $C_x$  centered at the point x, contained in  $\Omega$  and congruent to a finite cone of fixed aperture centered at the origin (a finite cone  $C_0$  is the intersection of open ball centered at the origin with the set  $\{\lambda x; \lambda > 0, |x - z_0| < r\}$ , where  $z_0 \neq 0$  and r > 0 are fixed), see for instance, [1], p. 300.

PROPOSITION 3.1. Let  $\Omega$  be an arbitrary domain in  $\mathbb{R}^n$ . Then for all  $x \in \Omega$ 

$$a_{\Omega}(x) \le \frac{c_1}{[\delta(x)]^{\alpha}}, \qquad c_1 = \frac{|S^{n-1}|}{\alpha} = \frac{2\pi^{\frac{1}{2}}}{\alpha\Gamma(\frac{n}{2})}.$$
 (3.1)

If the domain  $\mathbb{R}^n \setminus \overline{\Omega}$  has the cone property, then there exists a constant  $c_2 > 0$  such that

$$\frac{c_2}{[\delta(x)]^{\alpha}} \le a_{\Omega}(x) \le \frac{c_1}{[\delta(x)]^{\alpha}}.$$
(3.2)

P r o o f. For  $x \in \Omega$ , we have

$$\int_{\mathbb{R}^n \setminus \Omega} \frac{dy}{|x-y|^{n+\alpha}} \le \int_{\substack{y \in \mathbb{R}^n \\ |x-y| \ge \delta(x)}} \frac{dy}{|x-y|^{n+\alpha}} = \int_{\substack{y \in \mathbb{R}^n \\ |y| > \delta(x)}} \frac{dy}{|y|^{n+\alpha}}.$$
 (3.3)

Passing to polar coordinates, we arrive at the estimate in (3.1).

To prove the left-hand side estimate in (3.2), we choose the boundary point  $x_0 \in \partial \Omega$  (depending on x and not necessarily unique) at which  $|x - x_0| = \delta(x)$ . Then,

$$|x-y| \le |x-x_0| + |x_0-y|$$
$$a_{\Omega}(x) \ge \int_{R^n \setminus \Omega} \frac{dy}{[|x_0-y| + \delta(x)]^{n+\alpha}}.$$

Since  $\mathbb{R}^n \setminus \overline{\Omega}$  has the cone property, there exists a finite cone  $\Gamma_{\Omega}(x_0, \theta)$ with vertex at  $x_0$  and fixed aperture  $\theta (= \operatorname{arctg} \frac{r}{|z_0|}) > 0$ , such that  $\Gamma_{\Omega}(x_0, \theta) \subset \mathbb{R}^n \setminus \overline{\Omega}$ . Then

$$\int\limits_{R^n \setminus \Omega} \frac{dy}{[|x_0 - y| + \delta(x)]^{n + \alpha}} \ge \int\limits_{B(x_0, \delta(x)) \cap \Gamma_\Omega(x_0, \theta)} \frac{dy}{[|x_0 - y| + \delta(x)]^{n + \alpha}} \,.$$

After translation to the origin and passing to polar coordinates, we obtain  $\delta(m)$ 

$$a_{\Omega}(x) \geq \int_{0}^{\delta(x)} \frac{\rho^{n-1} d\rho}{[\rho+\delta(x)]^{n+\alpha}} \int_{S^{n-1}\cap\Gamma_{\Omega}(0,\theta)} d\sigma = \frac{c_2}{[\delta(x)]^{\alpha}},$$

where

$$c_2 = C(\Omega) \int_0^1 \frac{t^{n-1}dt}{(t+1)^{n+\alpha}} = \frac{C(\Omega)}{n} F(n, n+\alpha; n+1; -1), \quad C(\Omega) = |S^{n-1} \cap \Gamma_\theta(0)|$$

where  $F = {}_2F_1$  is the Gauss hypergeometric function.

4. The function 
$$a_{\Omega}(x)$$
 in the case of the ball

THEOREM 4.1. In case  $\Omega$  is the ball B(0, R), the function  $a_{\Omega}(x)$  has the form

$$a_{B(0,R)}(x) = \frac{|S_{n-1}|}{\alpha R^{\alpha}} F\left(\frac{\alpha}{2}, \frac{\alpha+n}{2}; \frac{n}{2}; \frac{|x|^2}{R^2}\right), \qquad |x| < R.$$
(4.1)

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and

It has also the following representation

$$a_{B(0,R)}(x) = \frac{\mathcal{A}_R(|x|)}{(R^2 - |x|^2)^{\alpha}}$$
(4.2)

with  $\mathcal{A}_R(0) = \frac{1}{R^{\alpha}} |S^{n-1}|$  and  $\mathcal{A}_R(R) = (2R)^{\alpha} \frac{\pi^{\frac{n-1}{2}} \Gamma(\frac{1+\alpha}{2})}{\alpha \Gamma(\frac{n+\alpha}{2})}$ , where

$$\mathcal{A}_{R}(r) = \frac{\pi^{\frac{n}{2}}}{\Gamma\left(1+\frac{\alpha}{2}\right)\Gamma\left(\frac{n-\alpha}{2}\right)} \int_{0}^{1} s^{\frac{\alpha}{2}-1} (1-s)^{\frac{n-\alpha}{2}-1} (R^{2}-r^{2}+sr^{2})^{\frac{\alpha}{2}} ds \quad (4.3)$$
$$= \frac{|S^{n-1}|}{\alpha} (R^{2}-|x|^{2})^{\frac{\alpha}{2}} F\left(-\frac{\alpha}{2},\frac{\alpha}{2};\frac{n}{2};-\frac{|x|^{2}}{R^{2}-|x|^{2}}\right).$$

P r o o f. The function  $a_{B(0,R)}(x)$  depends on |x| only. Let us denote  $a(r) := a_{B(0,R)}(x), r = |x|$ , for brevity. It suffices to consider the case R = 1. After passing to polar coordinates and using the Catalan formula

$$\int_{\substack{S^{n-1} \\ e \text{ at}}} f(x \cdot \sigma) \, d\sigma = |S^{n-2}| \int_{-1}^{1} f(|x|t)(1-t^2)^{\frac{n-3}{2}}, \quad x \in \mathbb{R}^n,$$

we arrive a

$$a(r) = \frac{|S_{n-2}|}{r^{\alpha}} \int_{-1}^{1} (1-t^2)^{\frac{n-3}{2}} dt \int_{\frac{1}{r}}^{\infty} \frac{s^{n-1}}{(s^2-2st+1)^{\frac{n+\alpha}{2}}} ds$$
$$= \frac{|S_{n-2}|}{r^{\alpha}} \int_{-1}^{1} (1-t^2)^{\frac{n-3}{2}} dt \int_{0}^{r} \frac{s^{\alpha-1}}{(s^2-2st+1)^{\frac{n+\alpha}{2}}} ds.$$

Since  $(1 - 2st + s^2)^{-\lambda} = \sum_{k=0}^{\infty} C_k^{\lambda}(t) s^k$ , where  $C_k^{\lambda}(t)$  are the Gegenbauer polynomials, we obtain

$$a(r) = 2|S_{n-2}| \sum_{k=0}^{\infty} \frac{r^{2k}}{2k+\alpha} \int_{0}^{1} (1-t^2)^{\frac{n-3}{2}} C_{2k}^{\frac{n+\alpha}{2}}(t) dt,$$
(4.4)

where we took into account that the Gegenbauer polynomials of odd order are odd. The formula

$$\int_{0}^{1} (1-t^{2})^{\frac{n-3}{2}} C_{2k}^{\frac{n+\alpha}{2}}(t) dt = \frac{\sqrt{\pi}\Gamma\left(\frac{n-1}{2}\right)}{2k!\Gamma\left(\frac{n}{2}\right)} \frac{\left(\frac{n+\alpha}{2}\right)_{k}\left(\frac{\alpha}{2}+1\right)_{k}}{\left(\frac{n}{2}\right)_{k}}$$

holds, see 2.21.2.3 in [3]. Making use of this formula, after easy calculations with the duplication formula for the Gamma-function taken into account, we arrive at

$$a(r) = \frac{1}{\alpha} |S_{n-1}| \sum_{k=0}^{\infty} \frac{\left(\frac{\alpha}{2}\right)_k \left(\frac{\alpha+n}{2}\right)_k}{k! \left(\frac{n}{2}\right)_k} r^{2k}$$

which is nothing else but (4.1) with R = 1.

To get (4.3) from (4.1), it suffices to make use of the transformation formula for the Gauss hypergeometric function:

$$F(a,b;c;z) = (1-z)^{-a} F\left(a,c-b;c;\frac{z}{z-1}\right).$$

COROLLARY 4.2. In the case n = 2m + 1 is odd,  $a_{B(0,R)}(x)$  is an elementary function for any  $\alpha \in (0,1)$ :

$$a_{B(0,1)}(x)\big|_{|x|=r} = d_{n,\alpha}(1-r^2)^{\frac{1-\alpha}{2}} \frac{d^m}{dr^{2m}} \left\{ (1-r^2)^{m-\frac{\alpha+1}{2}} \left[ (1+r)^{\alpha} + (1-r)^{\alpha} \right] \right\}$$
(4.5)

where  $d_{n,\alpha} = \frac{\alpha^{-1}\pi^{\frac{n-2}{2}}\Gamma(1-\frac{n-\alpha}{2})}{\Gamma(\frac{n+\alpha}{2})}$ . In particular, when n = 3 one has

$$a_{B(0,1)}(x) = \frac{2\pi}{\alpha^2 - 1} \left\{ \left( 2 + \frac{1}{\alpha} \right) \left[ (1 + |x|)^{-\alpha} + (1 - |x|)^{-\alpha} \right] + \frac{(1 - |x|)^{-\alpha} - (1 + |x|)^{-\alpha}}{|x|} \right\}.$$

P r o o f. When n is odd, we could write n = 2m + 1, then  $a(x) = \frac{1}{\alpha} |S_{n-1}| F\left(\frac{\alpha}{2}, \frac{\alpha+1}{2} + m; \frac{1}{2} + m; r^2\right).$ 

Using the formula

$$\frac{d^n}{dz^n} \left[ (1-z)^{a+n-1} F(a,b;c;z) \right]$$
  
=  $(-1)^n \frac{(a)_n (c-b)_n}{(c)_n} (1-z)^{a-1} F(a+n,b;c+n;z),$ 

see for example, 7.2.1.13 from [4], and the fact that the hypergeometric function is symmetric with respect to a and b, we get

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$$F\left(\frac{\alpha+1}{2}+m,\frac{\alpha}{2};\frac{1}{2}+m;r^{2}\right) = (-1)^{m}\frac{\left(\frac{1}{2}\right)_{m}}{\left(\frac{\alpha+1}{2}\right)_{m}\left(\frac{1-\alpha}{2}\right)_{m}}$$
$$\times (1-r^{2})^{\frac{1-\alpha}{2}}\frac{d^{m}}{dr^{2m}}\left[(1-r^{2})^{\frac{\alpha-1}{2}+m}F\left(\frac{\alpha+1}{2},\frac{\alpha}{2};\frac{1}{2};r^{2}\right)\right].$$
(4.6)

Then by the formula

$$F\left(a, a + \frac{1}{2}; \frac{1}{2}; z\right) = \frac{1}{2} \left[ (1 + \sqrt{z})^{-2a} + (1 - \sqrt{z})^{-2a} \right], \qquad (4.7)$$
formula 7.3.1.106, after simplifications we arrive at (4.5).

see [4], formula 7.3.1.106, after simplifications we arrive at (4.5).

REMARK 4.3. When n = 2m is even, the function a(r) may not be expressed in terms of elementary functions, being given by

$$a(r) = \frac{|S_{n-1}|}{\alpha} \cdot (-1)^{m-1} \frac{(m-1)!}{(\frac{\alpha}{2}+1)_{m-1}(1-\frac{\alpha}{2})_{m-1}} (1-r^2)^{-\frac{\alpha}{2}}$$
(4.8)  
 
$$\times \left(\frac{d}{dr^2}\right)^{m-1} \left\{ (1-r^2)^{m-2} \left[ P^0_{-\frac{\alpha}{2}-1} \left(\frac{1+r^2}{1-r^2}\right) - \frac{2r}{\alpha} P^1_{\frac{\alpha}{2}-1} \left(\frac{1+r^2}{1-r^2}\right) \right] \right\},$$

where  $P^0_{-\frac{\alpha}{2}-1} \equiv P_{-\frac{\alpha}{2}-1}$  and  $P^1_{-\frac{\alpha}{2}-1}$  are the Legendre polynomials and the associated Legendre function of the 1st kind. In particular, for n = 2 one has

$$a(r) = \frac{2\pi}{\alpha} (1 - r^2)^{-\frac{\alpha}{2} - 1} \left[ P^0_{-\frac{\alpha}{2} - 1} \left( \frac{1 + r^2}{1 - r^2} \right) - \frac{2r}{\alpha} P^1_{-\frac{\alpha}{2} - 1} \left( \frac{1 + r^2}{1 - r^2} \right) \right].$$
(4.9)

To obtain this, we use

$$F(a,b;a-b+2;z)$$

$$=\frac{\Gamma(a-b+2)}{b-1}z^{(b-a-1)/2}(1-z)^{-b}\left[aP_{-b}^{(b-a-1)/2}\left(\frac{1+z}{1-z}\right)-\sqrt{z}P_{-b}^{b-a}\left(\frac{1+z}{1-z}\right)\right],$$

see [4], formula 7.3.1.62, and (4.7) formula.

# 5. The operator $\mathbb{D}^{lpha}_{\Omega}$ as "quasi"- inverse to the operator $I^{lpha}_{\Omega}$

DEFINITION 5.1. The function  $\mu(x)$  is called a multiplier in the space X, if  $\mu f \in X$  and  $\|\mu f\|_X \leq c \|f\|_X$ , for all  $f \in X$ .

DEFINITION 5.2. Let  $\Omega$  be an open set in  $\mathbb{R}^n$ . We say that  $\Omega$  satisfies the *Strichartz condition* if there exist a coordinate system in  $\mathbb{R}^n$  and an integer N > 0 such that almost every line parallel to the axes intersects  $\Omega$ in at most N components.

The following statement shows that although the operator  $\mathbb{D}_{\Omega}^{\alpha}$  is not inverse to the operator  $I_{\Omega}^{\alpha}$  in the cases where  $\Omega \neq \mathbb{R}^{n}$ , it possesses some property of the inverse operator.

THEOREM 5.3. Let  $f = I_{\Omega}^{\alpha} \varphi$  where  $\varphi \in L_p(\Omega), 1 \leq p < \frac{1}{\alpha}$  and  $\Omega$  is a bounded domain satisfying the Strichartz condition. Then

$$\mathbb{D}_{\Omega}^{\alpha} f \in L_p(\Omega) \quad \text{and} \quad \|\mathbb{D}_{\Omega}^{\alpha} f\|_{L_p(\Omega)} \le C \|\varphi\|_{L_p(\Omega)}, \quad (5.1)$$

where C > 0 does not depend on f.

P r o o f. By the definition in (2.1) we have

$$\mathbb{D}^{\alpha}_{\Omega}I^{\alpha}_{\Omega}\varphi(x) = r_{\Omega}\mathbb{D}^{\alpha}\chi_{\Omega}I^{\alpha}\mathcal{E}_{\Omega}\varphi(x), \quad x \in \Omega.$$
(5.2)

Then the statement of the theorem is derived from the following three facts:

1) hypersingular integral operator is the left inverse operator to the Riesz potential operator in the case of the whole space  $\mathbb{R}^n$ :

$$\mathbb{D}^{\alpha} I^{\alpha} \varphi \equiv \varphi, \qquad \varphi \in L_p(\mathbb{R}^n), \qquad 1 \le p < n/\alpha,$$

see [5], p. 517.

2) the characteristic function  $\chi_{\Omega}$  of the domain  $\Omega$  satisfying the Strichartz condition is a multiplier in the space  $I^{\alpha}(L_p), L_p = L_p(\mathbb{R}^n)$  (see [6] for the case Bessel potentials and [2] for the case of Riesz potentials):

$$\|\mathbb{D}^{\alpha}\chi_{\Omega}I^{\alpha}\varphi\|_{L_{p}(\mathbb{R}^{n})} \leq C\|\varphi\|_{L_{p}(\mathbb{R}^{n})}, \qquad 1 
(5.3)$$

3) the condition  $\|\mathbb{D}^{\alpha} f\|_{L_{p}(\mathbb{R}^{n})} < \infty$  is sufficient for a function  $f \in L_{p}(\mathbb{R}^{n})$  to belong to  $I^{\alpha}(L_{p})$  and  $f = I^{\alpha}\mathbb{D}^{\alpha}f$ , see [5], Section 26. Observe that here we have used the fact that the domain  $\Omega$  is bounded: in case  $\Omega$  is unbounded, the function  $f = \mathcal{E}_{\Omega}I_{\Omega}^{\alpha}\varphi$  is not necessarily in  $L_{p}(\mathbb{R}^{n})$ .

Indeed, by (5.2) and (5.3) we have

$$\|\mathbb{D}^{\alpha}\mathcal{E}_{\Omega}I_{\Omega}^{\alpha}\varphi\|_{L_{p}(\mathbb{R}^{n})} = \|\mathbb{D}^{\alpha}\chi_{\Omega}I^{\alpha}\mathcal{E}_{\Omega}\varphi\|_{L_{p}(\mathbb{R}^{n})} \le C\|\mathcal{E}_{\Omega}\varphi\|_{L_{p}(\mathbb{R}^{n})} = C\|\varphi\|_{L_{p}(\Omega)}.$$

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Therefore, by 3) there exists a function  $\psi \in L_p(\mathbb{R}^n)$  such that  $\mathcal{E}_{\Omega}I^{\alpha}_{\Omega}\varphi(x) = I^{\alpha}\psi(x), x \in \mathbb{R}^n$  with  $\psi = \mathbb{D}^{\alpha}\mathcal{E}_{\Omega}I^{\alpha}_{\Omega}\varphi$ . Observe that  $\|\psi\|_{L_p(\mathbb{R}^n)} \leq C \|\varphi\|_{L_p(\Omega)}$  by (5.3). Then

$$\|\mathbb{D}^{\alpha}_{\Omega}I^{\alpha}_{\Omega}f\|_{L_{p}(\Omega)} = \|r_{\Omega}\mathbb{D}^{\alpha}\mathcal{E}_{\Omega}I^{\alpha}_{\Omega}\varphi\|_{L_{p}(\Omega)} = \|r_{\Omega}\mathbb{D}^{\alpha}I^{\alpha}\psi\|_{L_{p}(\Omega)}.$$

Consequently, by 1),

$$\|\mathbb{D}^{\alpha}_{\Omega}I^{\alpha}_{\Omega}\varphi\|_{L_{p}(\Omega)} = \|\psi\|_{L_{p}(\Omega)} \le \|\varphi\|_{L_{p}(\Omega)}$$

and (5.1) thus having been proved.

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